



A Design of a Low-Reynolds Number Airfoil that Leads to the Formation of Separation Bubbles at the Leading Edge

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Author's contribution

The sole author designed, analyzed and interpreted and prepared the manuscript.

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Abstract

The aerodynamics of airfoils at low Reynolds numbers (Re) has become increasingly important from both fundamental and industrial points of view, due to recent developments in small wind turbines, small-unmanned aerial vehicles (UAVs), micro-air vehicles (MAVs), as well as researches on bird/insect flying aerodynamics. Researchers on airfoil aerodynamics have focused on conventional aircraft design with Re beyond 5×10^5 , where separation bubble forming on the leading edge has positive impact on aerodynamic performance of the airfoil, increasing the lift. Separation bubble not form at low Re ($< 1.0 \times 10^4$), causing the value of lift to be small. The shape of the airfoil is change to get an airfoil that can cause separation at the leading edge. This is made possible by varying the maximum thickness, maximum camber and position of maximum camber to obtain different airfoils and their pressure difference computed. Reynolds numbers considered here are 1000 to 10, 000. It is from the study that a desired airfoil is obtained that has high lift and leading edge separation. In this study, it is now possible to design airfoil that can work well at low Reynolds number that is where velocity is low.

Keywords: Low-Reynolds number; leading edge separation; airfoil; lift and drag.

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Abbreviations

<i>LEV</i>	: <i>Leading edge vortex</i>
<i>Re</i>	: <i>Reynolds number</i>
<i>MAVS</i>	: <i>Micro-air vehicles</i>
<i>UAVS</i>	: <i>Unmanned aerial vehicles</i>
<i>LEB</i>	: <i>Leading edge separation</i>
<i>M</i>	: <i>Maximum camber,</i>
<i>P</i>	: <i>Maximum camber location,</i>
<i>t</i>	: <i>Maximum thickness</i>

1 Introduction

Rivers, tidal flows and wind are important in the search for alternative renewable energy source which when properly harnessed can have significant contribution towards meeting the increasing demand for clean energy.

Airfoils work well at high Reynolds number that is producing high lift that comes because of separation bubbles that form at the leading edge.

It is well known that the aerodynamic performance of airfoils that are optimal for high-speed therefore high Reynolds number significantly degrades when used for low-Reynolds-number applications where the Reynolds numbers are smaller.

[1] was perhaps the first to study separation bubbles? He observed the existence of separation and reattachment of the boundary layer over cambered airfoil.

Ever since the first observation by [1], the leading edge separation (LES) had been extensively investigated by numerous researchers, as reviewed by [2] over a low Reynolds number airfoil. If the free stream turbulence intensity is low, the flow starts as laminar; before transition, the laminar boundary layer separates due to the adverse pressure gradient. The separated flow quickly undergoes transition and becomes turbulent. Depending on parameters such as the local Reynolds number, pressure gradient, surface roughness, and free stream turbulence intensity, the turbulent free shear layer may entrain to reattach as a turbulent boundary layer behind a laminar separation bubble.

[2] Realized that steady-state aerodynamics does not accurately account for the forces produced by natural fliers, and this prompted several studies on the unsteady flow produced. Mechanisms such as rotational circulation, wake capture, and the unsteady leading edge vortex do account for the aerodynamics forces. Regarding forward flight, the unsteady leading edge vortex was the only mechanism present to produce the necessary forces. The unsteady leading edge vortex involves leading edge flow separation that reattaches to the wing and forms a separation bubble.

A prominent leading edge vortex (*LEV*), the hallmark of dynamic stall, has been observed on the leading edge of model Manduca wings at $Re=5000$ and model Drosophila wings at $Re=150$. In Drosophila, this enlarged area of vorticity is prominent at angles of attack above $\sim 12^\circ$, at which flow separates from the leading edge [2].

The vortex increases the circulation around the wing and creates much higher lift than the steady-state case. Within nature the primary unsteady a recirculation produced during the wing's downstroke. It acts to create a region of low pressure over the upper surface of the wing, although it can also be considered as augmenting the circulation around the wing, and thus increases lift. Aerodynamic phenomenon responsible for lift augmentation is the Leading Edge Vortex (*LEV*). The *LEV* is a region of highly three-dimensional.

There are three different types of separation depending on the position there situated; firstly: Trailing-edge stall, it is a separation that occur at the trailing edge and moving towards the leading edge as the incidence angle increases it occurs typically on thick airfoil. Secondly, Leading-edge stalls, caused by an abrupt separation of the flow near the leading edge without subsequent reattachment. The bubbles in this category are short; it also increases with angle of attack. Lastly, thin-aerofoil stall which is the flow separation at the leading edge with reattachment at a point by [3].

Studies have relied on mechanical wing flappers designed to mimic kinematics over a wide range of Reynolds numbers. Each of these studies has identified a leading-edge vortex (*LEV*) responsible for a significant portion of the lift generated. The strength and stability of this *LEV* depended on a number of variables, the Reynolds number, the wing shape, the translation, the rotation, and the angle of attack. The discovery of leading-edge vortices (*LEVs*) on the wings of insects in flight greatly advanced the knowledge of their dominant lift-generating mechanisms. Sharp leading edges induce high lift production through flow separation with vortical flow attached to the upper surface of insect wings during flapping and gliding.

The same studies were carried by [4] on the insect wings, most of the lift was associated with a large, stable leading edge vortex which separates from the sharp leading edge of the wing. Also model of swift wings done by [5] and another model on bats done by Edmonds, (2005) and [6] indicated a leading edge vortex lift enhancement mechanism at low Reynolds number regime which characterizes insects.

According to quasi-stead-state aerodynamic theory, slow-flying vertebrates are not able to generate enough lift to remain aloft. Therefore, unsteady aerodynamic mechanisms to enhance lift production was studied by [6] using digital particle image velocimetry and showed that small nectar-feeding bat is able to increase lift by as much as 40% using attached leading edge vortices (*LEVS*) during slow forward flight, resulting in a maximum lift coefficient of 4.8.

After focusing on a fluid dynamic around an airfoil and expressing the shape of airfoils as a function of three parameters, which control, the maximum thickness, maximum camber and maximum camber location of airfoil, a desired shape of airfoil that caused the formation of separation bubble at the leading edge at low Reynolds number can be found.

2 Literature Review

The aerodynamic design methods and principles developed over the past 40 years have produced efficient airfoils for conventional, large-scale, high-speed aircraft whose Reynolds numbers is beyond 1.0×10^5 . There is considerable literature on biological flight mechanisms; but there is very little detailed aerodynamic research available. Most hovering animals, such as insects and hummingbirds, enhance lift by producing leading edge vortices (*LEVs*) and by using both the downstroke and upstroke for lift production by [7] *LEVs* have been observed on butterfly wings in free flight [8]. Aerodynamics at the Reynolds numbers is considerably different from those of more conventional aircraft. The flow is laminar and viscously dominated. Boundary layers are quite thick, often reaching a significant fraction of the chord length.

The dependency of separation bubble on Reynolds number was first founded by [9] the study was made of laminar separation bubbles formed over a wide range of Reynolds numbers and in a variety of pressure distributions. His conclusion was that, the structure of the bubbles depended on the value of the Reynolds number of the separating boundary layer and a parameter based on the pressure rise over the region occupied by the bubbles. Conditions for the bursting of 'short' bubbles were determined by a unique relationship between these two parameters.

The tail (caudal fin) is one of the most prominent characteristics of fishes. Vortex reattachment was seen at the leading of the fish tail by [10] using 3-D high-resolution numerical of self-propelled virtual swimmers with different tail shapes. He showed that the evolution of the *LEV* drastically alters the pressure distribution on the tail and the force is generated.

Flow around the leading edge of the oscillating airfoil play a key role in the development and shedding of a concentrated vortex known as the dynamic stall vortex (cause of leading edge vortex). The favorable effect of the vortex was to increase lift by [11].

The importance of the *LEV* was noted by [11] and [12] in the context of Weis-Fogh's 'clap-and-fling' mechanism. The formation of an *LEV* was examined on both tethered and model dragonfly wings by Luttges et al. (1985). In a seminal study [13] visualized an *LEV* on the wing of a live hawk moth in tethered flight ($Re \sim 4000$). Efforts have been made to yield high lift flow mechanisms, [13,14] studied insect wings, they found out that most of the lift is associated with a large, stable leading edge vortex which separates from the sharp leading edge of the wing.

[15] Showed that the leading edge vortex is stable on wings steadily rotating like a propeller. The resultant aerodynamic force is normal to the wing surface, reflecting the fact that the leading edge vortex essentially eliminates the leading edge suction.

[16] studied the laminar separation bubble formed by an airfoil with cylindrical leading edge, followed by a constant thickness section, followed tapered trailing edge which formed laminar separation bubbles near the leading edge with Re ranging from 1.5×10^5 to 4.5×10^5 .

[17] and [18] investigated the incompressible and compressible flow around the leading edge; he found that the formation and location of separation bubbles depends on angle of attack and leading edge curvature.

A study by [19] showed that at the inner wing of flying hawkmoth there is a single attached LEV, while at mid wing there are multiple LEVs and that the strong and complex LEV suggests high flight power in hawkmoths

3 Methodology

3.1 Introduction

An Inviscid flow can be solved in different ways, such as the use of panel method. Under this method, we have variety depending with the choice of the singularity used to represent the velocity potential on the airfoil (sources, doublets or vortices).

The present work uses sources distribution; the method will provide tangential velocity distribution on the airfoil's surface (U_s). The pressure coefficient is then computed using the Bernoulli equation. The lift and drag coefficients are calculated by integrating the pressure coefficients over the body surface.

3.2 Computation of surface coordinates

The shape of an airfoil is expressed analytically as a function of three parameters which control, the maximum camber m , maximum camber location p , and maximum thickness t of the airfoil.

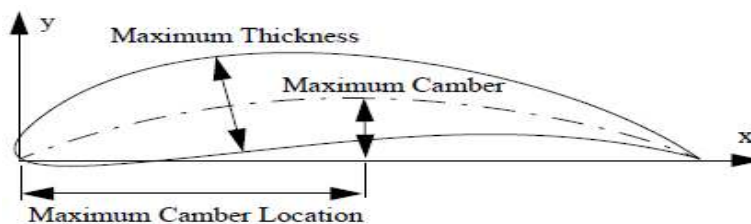


Fig. 1. Airfoil shape parameters

A wide variety of airfoils can be obtained by varying the three parameters.

Wing section is obtained by combining the camber line and the thickness distribution as shown.

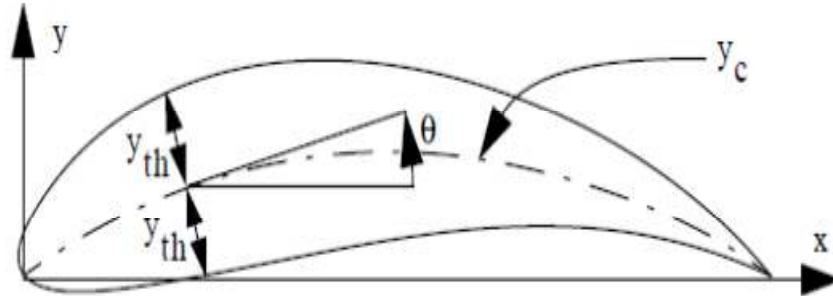


Fig. 2. Wing section

$$\left. \begin{aligned} x_u &= x - y_{th} \cos \Theta \\ x_l &= x + y_{th} \cos \Theta \end{aligned} \right\} \text{for } x < p \quad (1)$$

$$\left. \begin{aligned} x_u &= x + y_{th} \cos \Theta \\ x_l &= x - y_{th} \cos \Theta \end{aligned} \right\} \text{for } x > p \quad (2)$$

$$y_u = y_c + y_{th} \cos \Theta \quad (3)$$

$$y_l = y_c - y_{th} \cos \Theta \quad (4)$$

where (x_u, y_u) and (x_l, y_l) are points on the upper and lower surface respectively.

The thickness distribution and the camber line are given by;

$$y_{th} = 5tc \left(0.2969(x/c)^{1/2} - 0.126(x/c) - 0.3537(x/c)^2 + 0.2843(x/c)^3 + 0.1015(x/c)^4 \right) \quad (5)$$

$$\text{and } y_c = m/p^2 \left[2((x/c) - (x/c)^2) \right] \quad \text{for } (x/c) \leq p \quad (6)$$

$$= m/(1-p)^2 \left[1 - 2p + 2p(x/c) - (x/c)^2 \right] \quad \text{for } (x/c) \geq p \quad (1)$$

In these expressions, c is the airfoil chord length, m is the maximum camber, p is the maximum camber location, and t is the maximum thickness.

3.3 Computation of tangential velocity

The flow is simulated around the airfoil using MATLAB, the airfoil surface is divided into piecewise straight line segments called panel. Each panel is treated as a uniform source panel. Each panel is emitting a constant source of fluid along its length parallel to the normal vector of each panel.

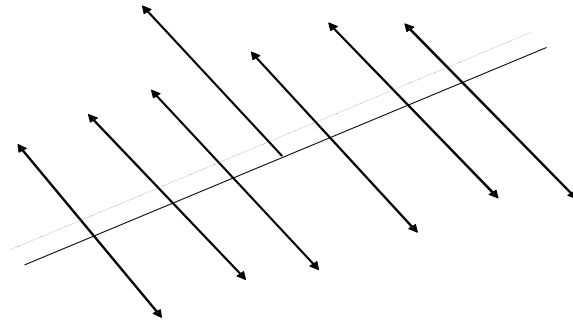


Fig. 3. A source panel

We define the normal vector to point outwards, towards the outside of the object as shown in Fig. 3.

The velocity of the flow in the radial direction is given by

$$U_r = \frac{m}{2\pi} \tag{8}$$

Where m is the source strength and determines the magnitude of the velocity, and r is the radial distance from the source. Since the airfoil was made up of multiple panels, the flow from each panel affects the flow at each other panel. Panels on the bottom of the airfoil induce a flow upwards on the top panels, and the top panels induce a flow downwards on the bottom panels.

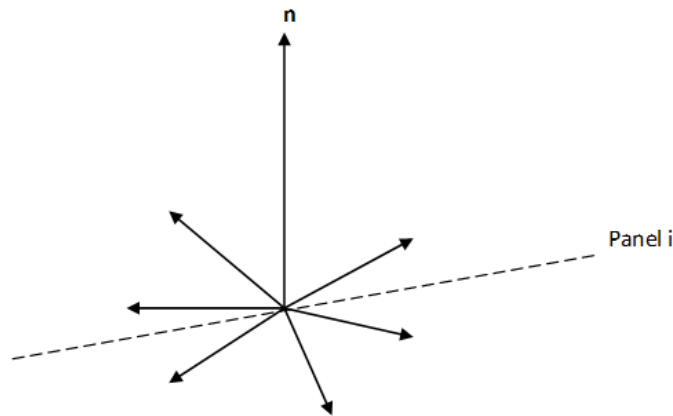


Fig. 4. Point source representation

To simplify things, each panel is treated as a point source rather than a source panel, as shown in Fig. 4. When a large number of panels are used, the size of each panel is small and they are represented as point sources without significant error.

We place the point source at the control point of each panel, which is located at the center. Finding the location of the control point equation (9) and (10) are used.

$$x_{cpi} = \frac{x_i + x_{i+1}}{2} \quad (9)$$

$$y_{cpi} = \frac{y_i + y_{i+1}}{2} \quad (10)$$

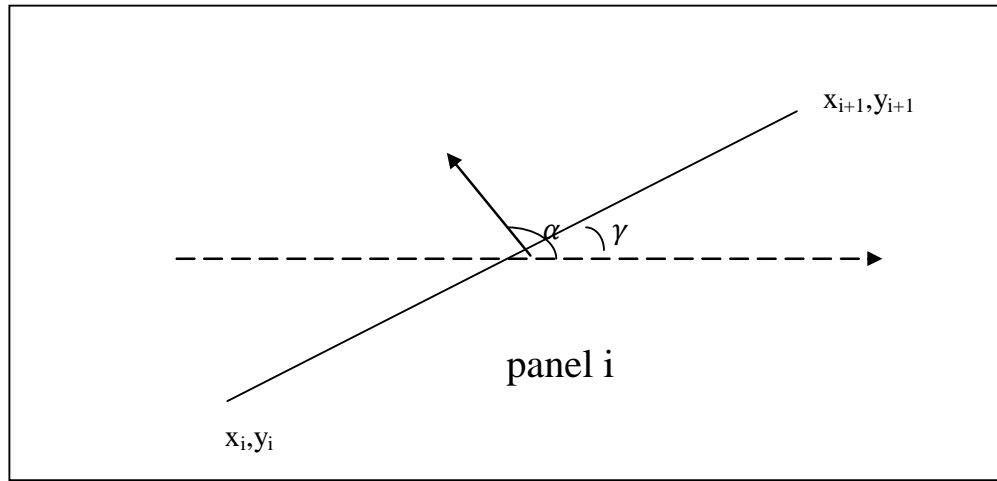


Fig. 5. Normal vector definition

If the end points of panel i are x_i, y_i and x_{i+1}, y_{i+1} as shown in Fig. 5, then the center point is located at x_{cpi}, y_{cpi} where x_{cpi} and y_{cpi} are given by Equations (9) and (10).

The normal vector at each control point is computed using two steps: 1) Find the angle between the normal vector and the positive x-axis. 2) Use sine and cosine to decompose the angle into a unit vector.

The angle γ , measured between the panel and the positive x-axis, is easily measured

Using arctangent

$$\gamma = \arctan \left(\frac{y_{i+1} - y_i}{x_{i+1} - x_i} \right) \quad (11)$$

Finding α simply required adding $\frac{\pi}{2}$ to γ ;

$$\alpha = \frac{\pi}{2} + \gamma = \frac{\pi}{2} + \arctan \left(\frac{y_{i+1} - y_i}{x_{i+1} - x_i} \right) \quad (12)$$

The normal vector is then defined as:

$$\mathbf{n} = (\cos \alpha \mathbf{i}, \sin \alpha \mathbf{j}) \quad (13)$$

Similarly, the tangent vector, is defined as:

$$\mathbf{t} = [\cos \left(\alpha - \frac{\pi}{2} \right) \mathbf{i} \sin \left(\alpha - \frac{\pi}{2} \right) \mathbf{j}] \quad (14)$$

Another value that is required is the length of the panel.

That is found using the following equation:

$$li = \sqrt{(x_{i+1} - x_i)^2 + (y_{i+1} - y_i)^2} \quad (15)$$

Where; li → the length of the panel

(x_i, y_i) and (x_{i+1}, y_{i+1}) → the end points of panel i

The velocity from one panel acting on another panel is considered

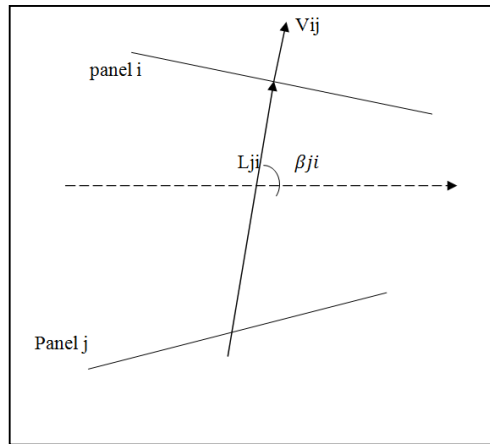


Fig. 6. Panel-to-panel induced velocity

As shown in Fig. 6, panel j induced a velocity, V_{ij} on panel i. The distance measured between the control points of panel j and panel i is represented by L_{ji} . The angle between the ray connecting panel j's control point to panel i's control point made an angle β_{ji} with the positive x-axis. Equations 16 and 17 are used to find the values the L_{ji} and β_{ji} :

$$L_{ji} = \sqrt{((x_{cpi} - x_{cpj})^2 + (y_{cpi} - y_{cpj})^2)} \quad (16)$$

$$\beta_{ji} = \arctan \left(\frac{y_{cpi} - y_{cpj}}{x_{cpi} - x_{cpj}} \right) \quad (17)$$

The vector V_{ij} was decomposed two into components; one normal to panel i and one tangent to the panel.

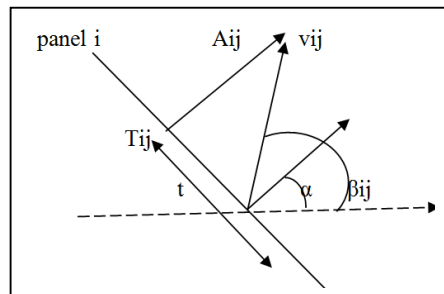


Fig. 7. Decomposing v_{ij}

The vector A_{ij} that represents the normal component of V_{ij} , and T_{ij} is the tangential component. Using equation 8, V_{ij} is written as:

$$V_{ij} = \frac{m_j}{2\pi L_{ji}} (\cos(\beta_{ji})\hat{i}, \sin(\beta_{ji})\hat{j}) \quad (18)$$

To get the normal component of V_{ij} , we take the dot product of velocity with normal vector.

$$\begin{aligned} V_{ijn} = V_{ij} \cdot \hat{n}_i &= \frac{m_j}{2\pi L_{ji}} (\cos(\beta_{ji})\hat{i} \sin(\beta_{ji})\hat{j}) \cdot (\cos\alpha\hat{i}, \sin\alpha\hat{j}) \\ &= \frac{m_j}{2\pi L_{ji}} (\cos\beta_{ji}\cos\alpha + \sin\beta_{ji}\sin\alpha) \end{aligned} \quad (2)$$

the normal influence Coefficient, A_{ij} is obtained by dividing normal component by m_j

$$A_{ij} = \frac{V_{ijn}}{m_j} = \frac{1}{2\pi L_{ji}} (\cos\beta_{ji}\cos\alpha + \sin\beta_{ji}\sin\alpha) \quad (20)$$

Applying a trigonometric identity, Equation is then written as:

$$A_{ij} = \frac{1}{2\pi L_{ji}} \cos(\beta_{ji} - \alpha) \quad (21)$$

Similarly, the tangential influence coefficient is written as the dot product of V_{ij} and \hat{t} , divided by m_j :

$$\begin{aligned} v_{ij,t} &= \frac{m_j}{2\pi L_{ji}} (\cos(\beta_{ji})\hat{i}, \sin(\beta_{ji})\hat{j}) \cdot [\cos(\alpha - \frac{\pi}{2})\hat{i} \sin(\alpha - \frac{\pi}{2})\hat{j}] \\ T_{ij} = \frac{v_{ij,t}}{m_j} &= \frac{1}{2\pi L_{ji}} [\cos\beta_{ji}\cos(\alpha - \frac{\pi}{2}) + \sin\beta_{ji}\sin(\alpha - \frac{\pi}{2})] \\ &= \frac{1}{2\pi L_{ji}} \cos(\beta_{ji} - \alpha + \frac{\pi}{2}) \end{aligned} \quad (22)$$

After finding the velocity induced on one panel by another, we sum up all of the velocities induced on each panel. The normal components of the velocity induced on panel i is written as:

$$U_{ni} = \sum_{j=1}^n m_j A_{ij} + U_{\infty n_i} \quad (23)$$

where $U_{\infty n_i}$ is the free stream velocity dotted with the normal vector of the panel which yielded the portion of the free stream velocity that is normal to the panel. Having the boundary conditions for a solid airfoil; that there cannot be any flow through a panel. Then equation 23 is equated to zero.

$$\begin{aligned} \sum_{j=1}^n m_j A_{ij} + U_{\infty n_i} &= 0 \\ \begin{bmatrix} m_1 \\ \vdots \\ m_n \end{bmatrix} \begin{bmatrix} A_{i,1} & A_{i,2} & \dots & A_{i,n} \end{bmatrix} &= -|U_{\infty n_i}| \end{aligned} \quad (24)$$

For the entire system (all the panels), more rows are added to the A matrix:

$$\begin{bmatrix} m_1 \\ \vdots \\ m_n \end{bmatrix} \begin{bmatrix} A_{1,1} & \dots & A_{1,n} \\ \vdots & \ddots & \vdots \\ A_{n,1} & \dots & A_{n,n} \end{bmatrix} = - \begin{bmatrix} U_{n1} \\ \vdots \\ U_{nn} \end{bmatrix} \quad (25)$$

For simplicity's sake, the Equation is written as

$$[m] [A] = -[bn] \quad (26)$$

To find the source strengths for each panel, Equation 26 is solved.

$$[m] = [A]^{-1} [-bn] \quad (27)$$

Although the velocities normal to each panel sum to zero, the tangential velocities do not. The total tangential velocity at panel i is given by:

$$U_{si} = \sum_{j=1}^n m_j T_{ij} + U_{\infty} \cdot t_i \quad (28)$$

The Equation is also written as a matrix equation:

$$[U_s] = [T] [m] + [bt] \quad (29)$$

All three matrices on the right side of Equation are known, then U_s is calculated using basic multiplication matrix.

3.4 Computation of lift and drag coefficient

Finally, the coefficients of pressure at each panel's control point are calculated using Bernoulli's equation.

$$C_{pi} = 1 - \left(\frac{U_{si}}{U_{\infty}} \right)^2 \quad (30)$$

Looking for the lift coefficient C_l and drag coefficient C_d , dimensionless coefficients are obtained by integrating C_p (Cox. Et.al)

$$C_l = \int_0^1 (c_{pl} - c_{pu}) d \frac{x}{c} \quad (31)$$

$$C_d = \int_{-\frac{t}{2}}^{\frac{t}{2}} (c_{pF} - c_{pA}) d \frac{x}{c} \quad (32)$$

4 Results and Discussion

Surface coordinates of four different airfoils are obtained by using equations above and then MATLAB to run the program.

Airfoil 1 is an airfoil with the following three parameters: maximum camber 0.09, maximum camber location 0.1 and maximum thickness 0.09.

Airfoil 2 is an airfoil with the following three parameters: maximum camber 0.7, maximum camber location 0.1 and maximum thickness 0.09.

Airfoil 3 is an airfoil with the following three parameters: maximum camber 0.09, maximum camber location 0.5 and maximum thickness as 0.09.

Airfoil 4 is an airfoil with the following three parameters: maximum camber 0.09, maximum camber location 0.1 and maximum thickness 0.3.

Table 1. A table showing upper surface coordinate of four different airfoils

Airfoil 1		Airfoil 2		Airfoil 3		Airfoil 4	
x_u	y_u	x_u	y_u	x_u	y_u	x_u	y_u
0	0	0	0	0	0	0	0
0.1112	0.1264	0.1118	0.7363	0.1013	0.0707	0.1114	0.2115
0.2234	0.1322	0.2313	0.7300	0.2136	0.1052	0.2262	0.2345
0.3357	0.1287	0.3501	0.6945	0.3280	0.1245	0.3411	0.2331
0.4477	0.1187	0.4659	0.6336	0.4428	0.1309	0.4551	0.2165
0.5593	0.1035	0.5783	0.5495	0.5570	0.1256	0.5679	0.1888
0.6704	0.0838	0.6874	0.4437	0.6702	0.1095	0.6790	0.1524
0.7809	0.0599	0.7939	0.3167	0.7819	0.0830	0.7883	0.1087
0.8908	0.0320	0.8980	0.1689	0.8919	0.0464	0.8954	0.0579

Table 2. A table showing lower surface coordinate of four different airfoils

Airfoil 1		Airfoil 2		Airfoil 3		Airfoil 4	
x_l	y_l	x_l	y_l	x_l	y_l	x_l	y_l
1.0000	0.0000	1.0000	0.0000	1.0000	0.0000	1.0000	0.0000
0.8869	0.0097	0.8798	0.1555	0.8858	0.0247	0.8824	-0.0162
0.7746	0.0180	0.7617	0.2893	0.7736	0.0415	0.7673	-0.0308
0.6630	0.0249	0.6459	0.4013	0.6631	0.0505	0.6543	-0.0438
0.5519	0.0304	0.5328	0.4918	0.5541	0.0522	0.5432	-0.0549
0.4412	0.0349	0.4229	0.5614	0.4461	0.0469	0.4338	-0.0629
0.3310	0.0392	0.3166	0.6114	0.3387	0.0355	0.3256	-0.0652
0.2210	0.0445	0.2132	0.6442	0.2308	0.0192	0.2183	-0.0578
0.1110	0.0535	0.1104	0.6634	0.1209	0.0005	0.1108	-0.0315

After obtaining the surface coordinates value, the values are run in MATLAB to obtain pressure coefficients using equation 30.

Table 3. Tables showing upper pressure coefficients and lower pressure coefficients of airfoil 1 at different Reynolds number

Re=1000		Re= 2000		Re=3000		Re=4000	
-147.3087	0.8787	-36.0772	0.9697	-15.4787	0.9865	-8.2693	0.9924
0.8833	0.9645	0.9708	0.9911	0.9870	0.9961	0.9927	0.9978
0.9333	0.6226	0.9833	0.9056	0.9926	0.9581	0.9958	0.9764
0.9219	0.6447	0.9805	0.9112	0.9913	0.9605	0.9951	0.9778
0.9299	0.6328	0.9825	0.9082	0.9922	0.9592	0.9956	0.9770
0.9468	0.5882	0.9867	0.8971	0.9941	0.9542	0.9967	0.9743
0.9710	0.5373	0.9927	0.8843	0.9968	0.9486	0.9982	0.9711
0.6561	0.9454	0.9140	0.9863	0.9618	0.9939	0.9785	0.9966
0.8518	-45.6183	0.9630	-10.6546	0.9835	-4.1798	0.9907	-1.9136

Cont'

Re=5000		Re= 6000		Re=7000		Re8000	
cpu	Cpl	Cpu	Cpl	Cpu	cpl	cpu	cpl
-4.9323	0.9951	-3.1197	0.9966	-2.0267	0.9975	-1.3173	0.9981
0.9953	0.9986	0.9968	0.9990	0.9976	0.9993	0.9982	0.9994
0.9973	0.9849	0.9981	0.9895	0.9986	0.9923	0.9990	0.9941
0.9969	0.9858	0.9978	0.9901	0.9984	0.9927	0.9988	0.9944
0.9972	0.9853	0.9981	0.9898	0.9986	0.9925	0.9989	0.9943
0.9979	0.9835	0.9985	0.9886	0.9989	0.9916	0.9992	0.9936
0.9988	0.9815	0.9992	0.9871	0.9994	0.9906	0.9995	0.9928
0.9862	0.9978	0.9904	0.9924	0.9930	0.9989	0.9946	0.9991
0.9941	-0.8647	0.9959	0.9960	0.9970	0.0486	0.9977	0.2716

Cont'

Re=9000		Re= 10000		Re=100000	
cpu	cpl	Cpu	cpl	cpu	cpl
-0.8310	0.9985	-0.4831	0.9988	0.9852	1.0000
0.9986	0.9996	0.9996	0.9996	1.0000	1.0000
0.9992	0.9953	0.9993	0.9962	1.0000	1.0000
0.9990	0.9956	0.9992	0.9964	1.0000	1.0000
0.9991	0.9955	0.9993	0.9963	1.0000	1.0000
0.9993	0.9949	0.9995	0.9959	1.0000	1.0000
0.9996	0.9943	0.9997	0.9954	1.0000	1.0000
0.9958	0.9993	0.9966	0.9995	1.0000	1.0000
0.9982	0.4245	0.9985	0.5338	1.0000	0.9953

Table 4. Tables showing upper pressure coefficients and lower pressure coefficients of airfoil 2, airfoil 3 and airfoil 4 at Reynolds number 4000

Airfoil 2		Airfoil 3		Airfoil 4	
cpu	cpl	cpu	cpl	cpu	cpl
-14.9455	0.9971	0.8894	0.9934	0.9886	-1.6144
0.9852	0.9983	0.9958	0.9951	0.9685	0.9996
0.9887	1.0000	0.9926	0.9897	0.9715	0.9991
0.9931	0.9845	0.9914	0.9902	0.9737	0.9976
0.9970	0.9818	0.9914	0.9902	0.9750	0.9961
0.9990	0.9800	0.9924	0.9896	0.9752	0.9948
0.9704	0.9814	0.9941	0.9880	0.9738	0.9940
0.9934	0.9901	0.9885	0.9836	0.9751	0.9947
0.9968	-15.0872	0.9926	0.9872	-1.4273	0.9960

The values obtained of pressure coefficient from MATLAB are used to determine lift coefficient and drag coefficient. Equation 31 and 32 are used to compute. The efficiency of an airfoil is based on ratio of lift to drag or lift coefficient to drag coefficient. The ratios are obtained as shown below.

The results obtained on Table 5 shows that lift coefficient increases with decrease in Reynolds number. Drag coefficient increases also with decrease in Reynolds. The performance of airfoil is based on the ratios of the two coefficients. Looking on the ratios we can say that performance of Airfoil 1 is good at low Reynolds number. On the same we can see that when the Reynolds number is raised to 100,000 the ratio reduces, justifying that it can only work well at low Reynolds number.

The results at Table 6 show performances of the other airfoils. Since the results were obtain at Reynolds number of 4000, we will compare the results obtained with those on Table 5 where the Reynolds number is 4000. The ratio of Airfoil 1 is 3.1464, Airfoil 2 is -0.1077, Airfoil 3 is 0.0549 and that Airfoil 4 is -0.0191. it is from the results that we can say that it's only Airfoil 1 that can perform well.

Table 5. Table showing lift coefficient, drag coefficient and ratio of lift coefficient to drag coefficient of airfoil 1

	Re=1000	Re=2000	Re=3000	Re=4000	Re=5000	Re=6000	Re=7000	Re=8000	Re=9000	Re=10000	Re=100000
Cl	100.4105	25.1026	11.1567	6.2758	4.0164	2.7891	2.0492	1.5688	1.2397	1.0041	0.0101
Cd	10.3139	3.6585	2.4260	1.9946	1.7949	1.6865	1.6211	1.5787	1.5496	1.5287	1.4409
Ratio; $\frac{cl}{cd}$	9.7355	6.8614	4.5988	3.1464	2.2377	1.6538	1.2641	0.9937	0.8000	0.6568	0.0070

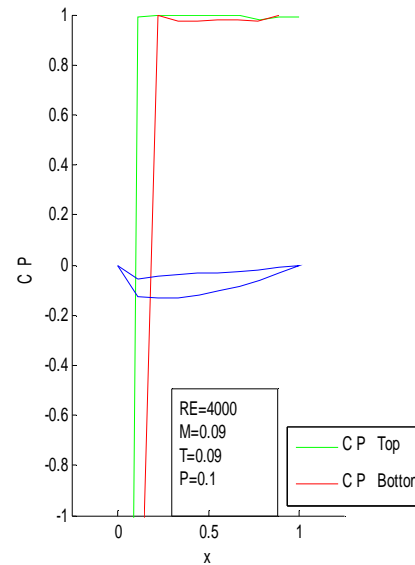


Fig. 8. Shape of Airfoil 1

4.1 Graphs showing the four shapes of airfoils

Graphs of pressure coefficient against chord length (x) are plotted from MATLAB. Also the shape of the airfoil is plotted corresponding to the parameters used. (Maximum camber, maximum camber location, and maximum thickness).

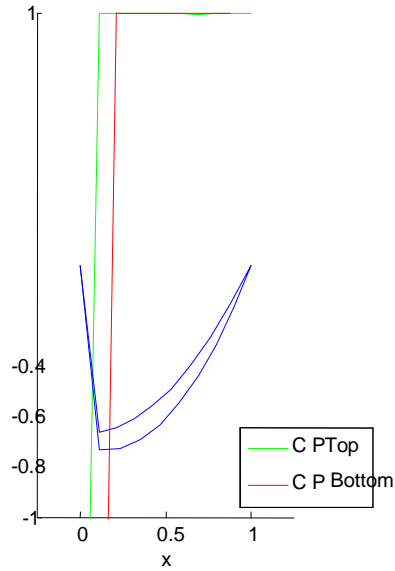


Fig. 9. Shape of airfoil 2

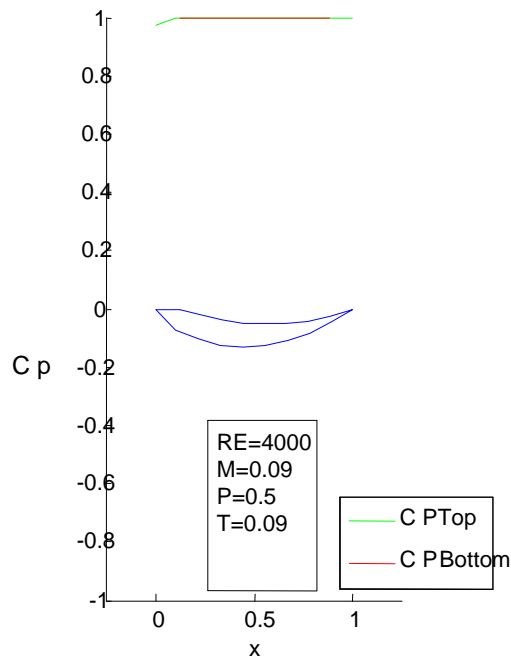


Fig. 10. Shape of airfoil 10

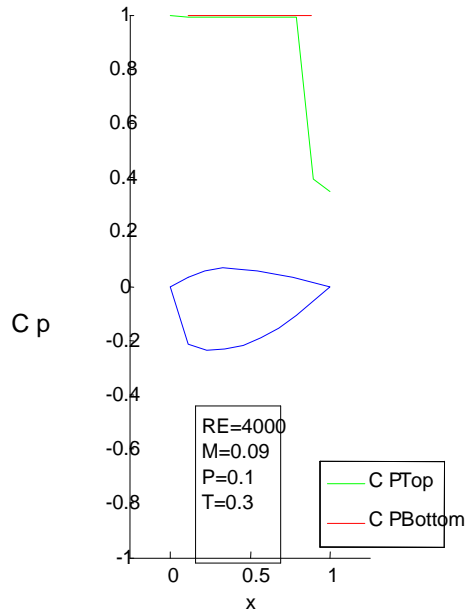


Fig. 11. Shape of airfoil 4

Table 6. Table showing lift coefficient, drag coefficient and ratio of lift coefficient to drag coefficient of airfoil 1, airfoil 2 and airfoil 3

	Airfoil 2	Airfoil 3	Airfoil 4
Cl	-0.1521	0.0788	-0.0166
Cd	1.4126	1.4361	0.8679
Ratio; $\frac{cl}{cd}$	-0.1077	0.0549	-0.0191

Looking at the shapes of the airfoil from the graphs we notice that they are different. From literature review we know that leading edge vortex causes separation bubbles at the leading edge. In the four airfoils it's Airfoil 1 that has a depression on the leading edge, so we expect the flow at that point to be vortical.

5 Conclusion and Recommendation

5.1 Conclusion

From the discussion on the lift coefficient on Table 5, we see that it is only one of the airfoils that is Airfoil 1 that has good performance in terms of lift coefficient (which is directly proportional to lift). On the same table we see the same airfoil operating with a different Reynolds number of 100000, and the lift coefficient at that point is seen to have dropped, implying that the airfoil can only work well at low Reynolds number.

We know that to get actual performance of the airfoil we need to check at the efficiency (the ratio of lift coefficient to drag coefficient or the ratio of lift to drag). Looking at the ratios, we see the same airfoil performing well at Reynolds number less than 10000 and beyond that its performance decreases. Also using a different airfoil the performance becomes poor.

Checking at the shape produced by Airfoil 1, it has a depression on the upper side on the leading edge. We expect the flow at the point to be vortical flow. It is from the vortical flow that we get the vortex, thus leading edge vortex (LEV) which causes separation at the leading edge. Therefore causing higher lift coefficient. The other airfoils lack the depression at the leading edge thus lacks the formation of the leading edge separation, hence low lift generation.

5.2 Recommendations

The research conducted so far has produced a desired airfoil that works well with low Reynolds number. I recommend this airfoil (turbine) be used to extract wind energy on areas where wind velocity is low, that is urban areas where obstruction from building, also in pacific Island countries where they have low speed of wind geographically.

6 Suggestions for Further Research

This research has focus in changing the shape of the airfoil to find an airfoil that can cause leading edge separation. It is suggested that further study be conducted to investigate how surface roughness, varying of the angle of attack can also cause the formation of leading edge separation at low Reynolds number of less than 10,000.

Competing Interests

Author has declared that no competing interests exist.

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Appendix

A program to compute surface coordinates and pressure coefficient of an airfoil using MATLAB.

```

clear;
clc;
close all;
%General variables
c = 1.000; %chord length
N = 5; %Positive integer number of points to use on one side
uinf = 1; %Strenght of the free stream velocity
aoa = 0*pi/180; %Angle between u inf and +x-axis in radians
%Converting to useful airfoil measurements
m = 60/100; %Maximum camber, in percent of chord
p = 8/10; %Position of max camber, in tenths of chord
t = 70/100; %Maximum thickness
%% Find the panel end points
% Define a vector of x coordinates
x = linspace(0, c, N); % Varied distances from 0 to chord along the airfoil
x = transpose(x);
%Calculate the height above mean chord line
yt = 5*t *c*(.2969*sqrt(x/c) - .126*(x/c) - .3516*(x/c).^2 + ...
.2843*(x/c).^3 - .1036*(x/c).^4);
%Calculate the mean camber line
yc = zeros(length(x), 1);
ycPrime = zeros(length(x), 1);
theta = zeros(length(x), 1);
for i = 1:length(x)
if(x(i)/c < p)
yc(i) = m *((2*p*x(i)/c)-(x(i)/c).^2)/p.^2;
ycPrime(i) = m *((2*p/c)-(2*x(i)/c).^2)/p.^2;
theta(i) = atan(m *((2*p/c)-(2*x(i)/c).^2)/p.^2);
else
yc(i) = m *(1-2*p+2*p*(x(i)/c)-(x(i)/c).^2)/(1-p).^2;
ycPrime(i) = m*(2*p/c-2*x(i)/c).^2/(1-p).^2;
theta(i) = atan(m*(2*p/c-2*x(i)/c).^2/(1-p).^2);
end
end
lowerPts = [x + yt.*sin(theta), yc - yt.*cos(theta)];
upperPts = [x - yt.*sin(theta), yc + yt.*cos(theta)];

%Make sure there is only one TE point and only one LE point
upperPts = upperPts(1:end-1, :); %delete the TE point
lowerPts = lowerPts(2:end, :); %delete the LE point
%Create a combination of all points in order moving CW around the airfoil
%starting at the LE
allPts = vertcat(upperPts, flipud(lowerPts));
%% Calculate various lengths, angles, and vectors
%Find the length of each panel
sideLen = zeros(length(allPts(:,1)),1);
for i = 1:length(allPts(:,1));
if(i + 1 > length(allPts))
sideLen(i) = sqrt((allPts(i, 1) - allPts(1, 1)).^2 + ...

```

```

(allPts(i, 2) - allPts(1, 2)).^2 );
else
sideLen(i) = sqrt((allPts(i, 1) - allPts(i + 1, 1)).^2 + ...
(allPts(i, 2) - allPts(i + 1, 2)).^2 );
end
end
%Find all control point locations and alpha values
CP = zeros(length(allPts(:,1)), 2);
alpha = zeros(length(allPts(:,1)), 1);
for i = 1:length(allPts(:,1))
if(i + 1 > length(allPts)) %connect the last point to the first point
alpha(i) = pi/2 + atan2(allPts(1, 2) - allPts(i, 2),...
allPts(1, 1) - allPts(i, 1));
CP(i,:) = allPts(i, :) + (allPts(1, :) - allPts(i, :))/2;
else
alpha(i) = pi/2 + atan2(allPts(i + 1, 2) - allPts(i, 2),...
allPts(i + 1, 1) - allPts(i, 1));
CP(i,:) = allPts(i, :) + (allPts(i + 1, :) - allPts(i, :))/2;
end
end
%Find the normal and tangent vectors for each panel
nVec = [cos(alpha), sin(alpha)];
tVec = [cos(alpha - pi/2), sin(alpha - pi/2)];

% Find distance between control points and the angle between each ray and
the free stream velocity.
L = zeros(length(CP(:,1)), length(CP(:,1)));
beta = zeros(length(CP(:,1)), length(CP(:,1)));
for i = 1:length(CP(:,1))
for j = 1:length(CP(:,1))
if(i == j)
L(i,j) = 0;
beta(i,j) = alpha(i);
else
L(i, j) = sqrt((CP(j, 1) - CP(i, 1)).^2 + (CP(j, 2) - CP(i, 2)).^2);
beta(i, j) = atan2(CP(j, 2) - CP(i, 2), CP(j, 1) - CP(i, 1));
end
end
end
%convert u to a vector
uvector = zeros(length(allPts(:,1)),2);
uvector(:,1) = cos(aoa);
uvector(:,2) = sin(aoa);
% Find the normal and tangential components of u wrt each panel
bn = zeros(size(nVec(:,1)));
bt = zeros(size(tVec(:,1)));
for i = 1:length(allPts(:,1))
bn(i) = -1*dot( uvector(i,:), nVec(i,:));
bt(i) = dot( uvector(i,:), tVec(i,:));
end
% Find the normal influence components and the tangential influence ...
normCoeff = zeros(length(CP(:,1)), length(CP(:,1)));
tanCoeff = zeros(length(CP(:,1)), length(CP(:,1)));
for i = 1:length(CP(:,1));

```

```

for j = 1:length(CP(:,1));
if(i == j)
normCoeff(i, j) = 1/(2*sideLen(i));
tanCoeff(i, j) = 0;
else
normCoeff(i, j) = cos(beta(j,i) - alpha(i))/(2*pi*L(j,i));

tanCoeff(i, j) = cos(beta(j,i) - alpha(i) + pi/2)/(2*pi*L(j,i));
end
end
end

% Use linear algebra to solve for the source strength of each panel
m = normCoeff\bn;
%Use the tangential influence coefficients to find the tangential
% velocities at each panel
vSi = tanCoeff*m + bt;
%Compute the coefficient of pressure
CPressure = 1 - (vSi/uinf).^2;

%Plot the points
plot(allPts(:, 1), allPts(:,2), ' *');
% Plot the airfoil and the coefficient of pressure
figure();
hold on;
plot(allPts(1:N, 1), CPressure(1:N), 'g');
plot(allPts(N+1:end, 1), CPressure(N+1:end), 'r');
plot(allPts(:, 1), -1*allPts(:,2));
%connect the last point to the first point
plot([allPts(end,1), allPts(1,1)], -1*[allPts(end,2), allPts(1,2)]);
axis([-0.25, 1.25, -1, 1]); %Set the xMin, xMax, yMin, yMax respectively
set(gca,'DataAspectRatio', [1 1 1]);
legend('C P _ Top', 'C P _ Bottom', 'Location', 'SE');
xlabel('x');
ylabel('C_ P');
set(gca, 'YDir', 'reverse');

```

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