

FE Dynamic Analysis Using Moving Support Element on Multi-Span Beams Subjected to Support Motions

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Abstract

In the present study, finite element dynamic analysis or time history analysis of two-span beams subjected to asynchronous multi-support motions is carried out by using the moving support finite element. The elemental equation of the element is based on total displacements and is derived under the concept of the quasi-static displacement decomposition. The use of moving support element shows that the element is very simple and convenient to represent continuous beam moving, deforming and vibrating simultaneously due to support motions. The comparison between the numerical results and analytical solutions indicates that the FE result agrees with the analytical solution.

Keywords

Moving Support Element, Support Motions, Rayleigh-Damped Bernoulli-Euler Beam, Multi-Span Beam, Time History Analysis, Finite Element Method

1. Introduction

Long and slender structures are often excited dynamically through support motions rather than by applied external loadings, e.g., piers, chimneys, towers [1], long bridges [2]-[9] and oil pipeline subjected to ground motions. These structures in turn are responding to support motion inputs. Due to the special feature of the excitation, the effective inertial loadings are applied on structures. The structures can be represented as Euler-Bernoulli beam subjected to multi-support motions if the foundation or ground soil is assumed to be rigid. Moreover, extended structures such as the Golden Gate Bridge and oil piping experience different ground motion at each support during an earthquake because the arrival time of seismic wave at each support is different. The vibration

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of the aforementioned structures is characterized as the problem of flexural vibration of beams with time-dependent boundary conditions. Mindlin and Goodman [10] developed the quasi-static decomposition method and applied it to obtain a solution of the problem. With this method, many researchers have investigated the structural response subjected to multiple support excitations by employing various techniques such as time history analysis, response spectrum method of analysis, frequency domain spectral analysis, etc. [11].

According to the quasi-static decomposition method, the transverse displacement of the beam subjected to support motions is composed of the quasi-static part and dynamic part. Statically determinate beams subjected to ground motions at supports are accompanied only by quasi-static displacement of rigid-body motion. Kim and Jhung [12] presented beam elements for statically determinate beams excited by support motions and showed the FE results agree with analytical solutions. However, statically indeterminate beams subjected to non-synchronous support motions are involved with not only the quasi-static displacement of rigid-body motion but also the quasi-static displacement of forced deformation. For the dynamic analysis of such beams, Kim [13] developed moving support element that can describe both static forced deformation and dynamic displacement. The author illustrated a single span fixed-hinged beam subjected to asynchronous support motions to show the element's performance. In this paper, two-span beams subjected to asynchronous multi-support motions are illustrated to show that the moving support element produces accurate dynamic responses even for the continuous beams. Since it is hardly possible to find the literature that compares numerical solutions with analytic ones, the numerical results including bending moment and shear force are compared with analytic solutions to show the high accuracy of the numerical results.

2. Multi-Span Beam Subjected to Support Motions

2.1. Rayleigh-Damped Euler-Bernoulli Beam

For a beam in flexure shown in **Figure 1(a)**, the transverse displacement at any point x and time t is denoted by $y(x, t)$ and the transverse force per unit length by $f(x, t)$. The system parameters are the mass per unit length $m(x)$ and the flexural rigidity $EI(x)$, where E is Young's modulus of elasticity and $I(x)$ is the cross-sectional area moment of inertia about an axis normal to x and y and passing through the center of the cross-sectional area. **Figure 1(b)** shows the free-body diagram corresponding to a beam element of length dx , where $V(x, t)$ denotes the shearing force and $M(x, t)$ the bending moment. According to simple beam theory, they are expressed as follows:

$$M(x, t) = EI(x) \frac{\partial^2 y(x, t)}{\partial x^2} \tag{1}$$

$$V(x, t) = -\frac{\partial}{\partial x} \left(EI(x) \frac{\partial^2 y(x, t)}{\partial x^2} \right) \tag{2}$$

The motion of a Rayleigh-damped Euler-Bernoulli beam with uniform cross-section is described by the following partial differential equation.

$$m\ddot{y}(x, t) + \left(2\alpha m + \beta EI \frac{\partial^4}{\partial x^4} \right) \dot{y}(x, t) + EI \frac{\partial^4 y(x, t)}{\partial x^4} = f(x, t) \tag{3}$$

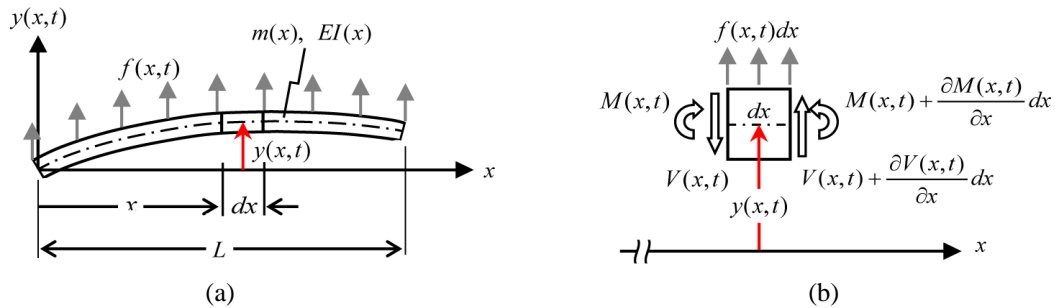


Figure 1. (a) Bending of a beam; (b) Free-body diagram of a beam element of length dx .

where a superimposed dot denotes a time derivative, L denotes length of the beam, and α and β are coefficients of the Rayleigh damping

2.2. Continuous Beam Subjected to Support Motions

For simplicity, we consider the dynamic response of two-span Rayleigh-damped Euler-Bernoulli beams subjected to multi-support excitation, which are shown in **Figure 2**, and assume that the external load $f(x, t)$ in Equation (3) is zero and that no other external loads are applied. The support motions of the beam are:

$$y(x, t)|_{x=0} = a_1(t) \quad (4)$$

$$y(x, t)|_{x=D_1} = a_2(t) \quad (5)$$

$$y(x, t)|_{x=L} = a_3(t) \quad (6)$$

where $a_i(t)$ ($i=1, 2$ and 3) are prescribed support displacements. Assume that the initial conditions are:

$$y(x, t)|_{t=0} = y_0(x) \quad \text{and} \quad \dot{y}(x, t)|_{t=0} = \dot{y}_0(x) \quad (7)$$

3. Moving Support Element

3.1. F.E. Equation Based on the Total Displacements

The moving support elemental equation is given, from [13], by

$$\left[m^{(e)} \right] \{ \ddot{y}^{(e)} \} + \left(2\alpha \left[m^{(e)} \right] + \beta \left[k^{(e)} \right] \right) \{ \dot{y}^{(e)} \} + \left[k^{(e)} \right] \{ y^{(e)} \} = \{ \hat{f}^{(e)} \} + \left[k^{(e)} \right] \left(\beta \{ \dot{y}_s^{(e)} \} + \{ y_s^{(e)} \} \right) \quad (8)$$

where

$$\left[m^{(e)} \right] = \int_0^l m \left[N \right]^T \left[N \right] d\hat{x} = \frac{ml}{420} \begin{bmatrix} 156 & 22l & 54 & -13l \\ 22l & 4l^2 & 13l & -3l^2 \\ 54 & 13l & 156 & -22l \\ -13l & -3l^2 & -22l & 4l^2 \end{bmatrix}, \quad (9)$$

$$\left[k^{(e)} \right] = \int_0^l EI \left[N'' \right]^T \left[N'' \right] d\hat{x} = \frac{EI}{l^3} \begin{bmatrix} 12 & 6l & -12 & 6l \\ 6l & 4l^2 & -6l & 2l^2 \\ -12 & -6l & 12 & -6l \\ 6l & 2l^2 & -6l & 4l^2 \end{bmatrix}. \quad (10)$$

In Equations (9) and (10), the double prime denotes a twice spatial differentiation with respect to the element coordinates \hat{x} depicted in **Figure 3** and the shape functions are:

$$\left[N \right] = \left[N_1(\hat{x}) \quad N_2(\hat{x}) \quad N_3(\hat{x}) \quad N_4(\hat{x}) \right] \quad (11)$$

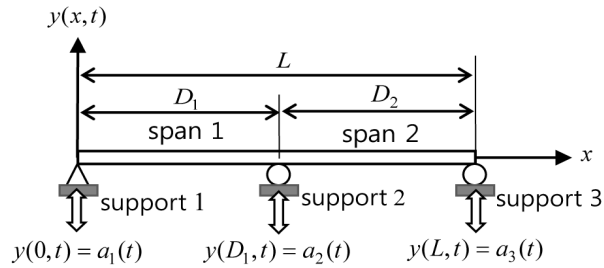


Figure 2. A two-span beam subjected to support motions, where $a_i(t)$ $i=1, 2, 3$ denote the support motions or the displacement time histories.

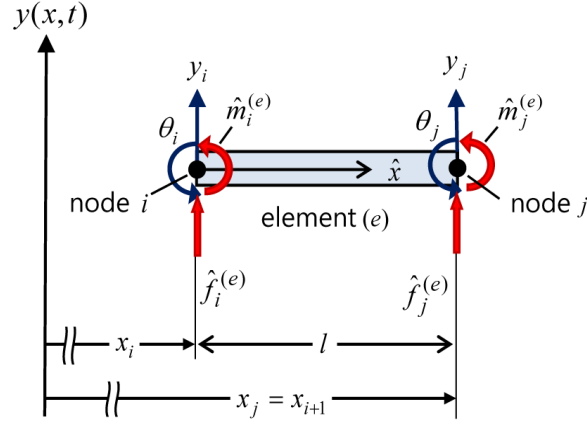


Figure 3. Typical beam element (e) subjected to concentrated nodal forces ($\hat{f}_i^{(e)}$ and $\hat{f}_j^{(e)}$) and moments ($\hat{m}_i^{(e)}$ and $\hat{m}_j^{(e)}$).

$$\begin{aligned} N_1(\hat{x}) &= \frac{1}{l^3}(2\hat{x}^3 - 3l\hat{x}^2 + l^3), & N_2(\hat{x}) &= \frac{1}{l^3}(l\hat{x}^3 - 2l^2\hat{x}^2 + l^3\hat{x}), \\ N_3(\hat{x}) &= \frac{1}{l^3}(-2\hat{x}^3 + 3l\hat{x}^2), & N_4(\hat{x}) &= \frac{1}{l^3}(l\hat{x}^3 - l^2\hat{x}^2) \end{aligned} \quad (12)$$

where l is element length. In Equation (8), $\{y^{(e)}\}$, $\{\dot{y}^{(e)}\}$ and $\{\ddot{y}^{(e)}\}$ are displacement, velocity and acceleration vector of element (e), respectively and they are:

$$\{y^{(e)}\} = \begin{Bmatrix} y_i \\ \theta_i \\ y_j \\ \theta_j \end{Bmatrix}, \quad \{\dot{y}^{(e)}\} = \begin{Bmatrix} \dot{y}_i \\ \dot{\theta}_i \\ \dot{y}_j \\ \dot{\theta}_j \end{Bmatrix}, \quad \{\ddot{y}^{(e)}\} = \begin{Bmatrix} \ddot{y}_i \\ \ddot{\theta}_i \\ \ddot{y}_j \\ \ddot{\theta}_j \end{Bmatrix} \quad (13)$$

where y_i and θ_i are transverse displacement and angular displacement at node i , respectively. The vectors $\{y_s^{(e)}\}$, $\{\dot{y}_s^{(e)}\}$ and $\{\ddot{y}_s^{(e)}\}$ are quasi-static displacement, velocity and acceleration vector of element (e), respectively and they are:

$$\{y_s^{(e)}\} = \begin{Bmatrix} y_{si} \\ \theta_{si} \\ y_{sj} \\ \theta_{sj} \end{Bmatrix}, \quad \{\dot{y}_s^{(e)}\} = \begin{Bmatrix} \dot{y}_{si} \\ \dot{\theta}_{si} \\ \dot{y}_{sj} \\ \dot{\theta}_{sj} \end{Bmatrix}, \quad \{\ddot{y}_s^{(e)}\} = \begin{Bmatrix} \ddot{y}_{si} \\ \ddot{\theta}_{si} \\ \ddot{y}_{sj} \\ \ddot{\theta}_{sj} \end{Bmatrix}. \quad (14)$$

where y_{si} and θ_{si} are quasi-static transverse displacement and angular displacement at node i , respectively. $\{\hat{f}^{(e)}\}^T = \{\hat{f}_i^{(e)} \quad \hat{m}_i^{(e)} \quad \hat{f}_j^{(e)} \quad \hat{m}_j^{(e)}\}$ is a local force vector of element (e).

Note that the underlined terms in right hand side of Equation (8) are peculiar to the moving support element and they contain the quasi-static displacement and velocity. The static components can be obtained exactly by static FE analysis, which will be considered in the next section.

3.2. Static FE Analysis for Quasi-Static Displacements

According to the quasi-static decomposition method, the solution can be decomposed into two parts:

$$y(x,t) = y_s(x,t) + w(x,t) \quad (15)$$

where y_s denotes the quasi-static displacement, and $w(x,t)$ is the dynamic contribution due to the inertial and damping effect. Using Equation (15), the total angular displacement $\theta(x,t) = \partial y(x,t)/\partial x$ is expressed as

follows:

$$\theta(x, t) = \theta_s(x, t) + \phi(x, t) \quad (16)$$

where $\theta_s(x, t)$ is quasi-static angular displacement and $\phi(x, t)$ is dynamic angular displacement. The quasi-static displacement $y_s(x, t)$ must satisfy

$$EI \frac{\partial^4 y_s(x, t)}{\partial x^4} = 0 \quad (17)$$

and Equation (17) is subjected to the support conditions in Equations (4)-(6).

Let Y_{sk} and Θ_{sk} ($k=1,2,3$) be quasi-static displacements and quasi-static angular displacements at the k -th support as shown in **Figure 4**, *i.e.*

$$Y_{s1} = y_s(x, t)|_{x=0}, \quad Y_{s2} = y_s(x, t)|_{x=D_1}, \quad Y_{s3} = y_s(x, t)|_{x=L}, \quad (18)$$

$$\Theta_{s1} = \theta_s(x, t)|_{x=0}, \quad \Theta_{s2} = \theta_s(x, t)|_{x=D_1}, \quad \Theta_{s3} = \theta_s(x, t)|_{x=L}. \quad (19)$$

The variables in Equations (18) and (19) will be called ‘quasi-static support variables’ or simply ‘support variables’ in this paper. The unknown support variables are determined by using the conventional static finite element method. The finite element equation is given by

$$\begin{bmatrix} 12 & 6D_1 & -12 & 6D_1 & 0 & 0 \\ 6D_1 & 4D_1^2 & -6D_1 & 2D_1^2 & 0 & 0 \\ -12 & -6D_1 & 24 & -6D_1 + 6D_2 & -12 & 6D_2 \\ 6D_1 & 2D_1^2 & -6D_1 + 6D_2 & 4D_1^2 + 4D_2^2 & -6D_2 & 2D_2^2 \\ 0 & 0 & -12 & -6D_2 & 12 & -6D_2 \\ 0 & 0 & 6D_2 & 2D_2^2 & -6D_2 & 4D_2^2 \end{bmatrix} \begin{Bmatrix} Y_{s1} \\ \Theta_{s1} \\ Y_{s2} \\ \Theta_{s2} \\ Y_{s3} \\ \Theta_{s3} \end{Bmatrix} = \begin{Bmatrix} F_{s1} \\ M_{s1} \\ F_{s2} \\ M_{s2} \\ F_{s3} \\ M_{s3} \end{Bmatrix}. \quad (20)$$

For the hinged beam in **Figure 2**, the support variables are given by

$$Y_{s1} = a_1(t), \quad Y_{s2} = a_2(t), \quad Y_{s3} = a_3(t) \quad (21)$$

and the external moments at supports are

$$M_{s1} = 0, \quad M_{s2} = 0, \quad M_{s3} = 0. \quad (22)$$

Using Equations (21) and (22), we obtain

$$\begin{aligned} \Theta_{s1}(t) &= A_{s1}^1 a_1(t) + A_{s1}^2 a_2(t) + A_{s1}^3 a_3(t), \\ \Theta_{s2}(t) &= A_{s2}^1 a_1(t) + A_{s2}^2 a_2(t) + A_{s2}^3 a_3(t), \\ \Theta_{s3}(t) &= A_{s3}^1 a_1(t) + A_{s3}^2 a_2(t) + A_{s3}^3 a_3(t) \end{aligned} \quad (23)$$

where A_{si}^k ($i=1,2,3; k=1,2,3$) are constants and they depends only on the span lengths. Note that the support variables for other beams can be determined by static FE method in the similar manner.

Using the support variables in Equations (21), (22) and (23), we obtain the distribution of the static displacement $y_s(x, t)$ and $\theta_s(x, t) = \partial y_s(x, t) / \partial x$:

$$y_s(x, t) = \begin{cases} N_1(\eta_1) \cdot Y_{s1}(t) + N_2(\eta_1) \cdot \Theta_{s1}(t) + N_3(\eta_1) \cdot Y_{s2}(t) + N_4(\eta_1) \cdot \Theta_{s2}(t) & \text{for } 0 \leq \eta_1 \leq D_1 \\ N_1(\eta_2) \cdot Y_{s2}(t) + N_2(\eta_2) \cdot \Theta_{s2}(t) + N_3(\eta_2) \cdot Y_{s3}(t) + N_4(\eta_2) \cdot \Theta_{s3}(t) & \text{for } 0 \leq \eta_2 \leq D_2 \end{cases} \quad (24)$$

$$\theta_s(x, t) = \begin{cases} \frac{dN_1(\eta_1)}{d\eta_1} \cdot Y_{s1}(t) + \frac{dN_2(\eta_1)}{d\eta_1} \cdot \Theta_{s1}(t) + \frac{dN_3(\eta_1)}{d\eta_1} \cdot Y_{s2}(t) + \frac{dN_4(\eta_1)}{d\eta_1} \cdot \Theta_{s2}(t) & \text{for } 0 \leq \eta_1 \leq D_1 \\ \frac{dN_1(\eta_2)}{d\eta_2} \cdot Y_{s2}(t) + \frac{dN_2(\eta_2)}{d\eta_2} \cdot \Theta_{s2}(t) + \frac{dN_3(\eta_2)}{d\eta_2} \cdot Y_{s3}(t) + \frac{dN_4(\eta_2)}{d\eta_2} \cdot \Theta_{s3}(t) & \text{for } 0 \leq \eta_2 \leq D_2 \end{cases} \quad (25)$$

where η_1 and η_2 are span coordinates and they are $\eta_1 = x$ and $\eta_2 = x - D_1$ as shown in **Figure 4**. Note

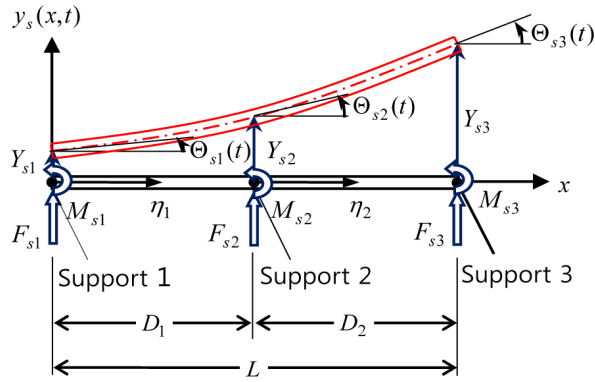


Figure 4. $Y_{s1}, Y_{s2}, Y_{s3}, \Theta_{s1}, \Theta_{s2}$ and Θ_{s3} are support variables and $F_{s1}, F_{s2}, F_{s3}, M_{s1}, M_{s2}$ and M_{s3} are external forces and moments applied at supports.

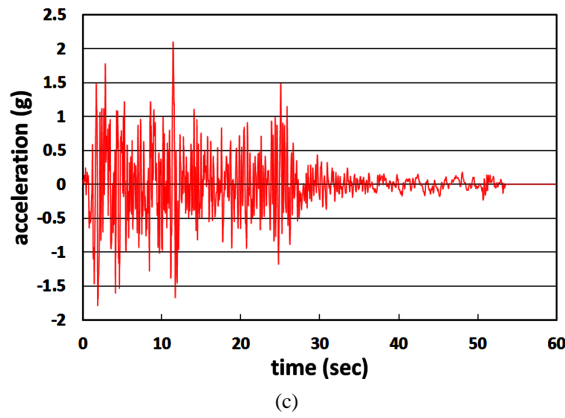
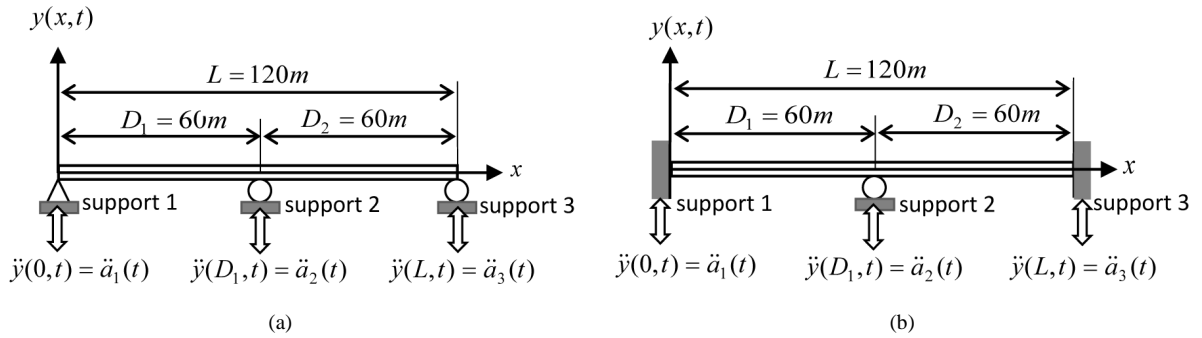


Figure 5. The beams in (a) and (b) are subjected to the seismic acceleration time histories when the earthquake traveling wave propagates longitudinally from the left support to the right ones at constant speed and the time delay between neighboring supports is 0.1 s. The seismic acceleration time history applied on the support 1 is shown in (c).

that the displacements $y_s(x,t)$ and $\theta_s(x,t)$ are exact ones. It is noteworthy that the quasi-static linear motion ($y_s(x,t), \dot{y}_s(x,t)$ and $\ddot{y}_s(x,t)$) and the quasi-static angular motion ($\theta_s(x,t), \dot{\theta}_s(x,t)$ and $\ddot{\theta}_s(x,t)$) at every time step can be determined exactly by static finite element analysis. Then using Equations (24) and (25), the vectors in Equation (14) are determined easily.

4. Numerical Tests

The two beams in **Figure 5(a)** and **Figure 5(b)** will be tested and the comparison between finite element outputs

and analytical solutions will be made to check the validity of the moving support element for dynamic responses of the beams due to support motions.

The input data are as follows: $D_1 = D_2 = 60$ m, $EI = 2.45 \times 10^9$ N·m², and $m = 2400$ kg/m; $\alpha = 0.0844$ s⁻¹ and $\beta = 0.0141$ s for the beam in **Figure 5(a)**; $\alpha = 0.1281$ s⁻¹ and $\beta = 0.0094$ s for the beam in **Figure 5(b)**. The initial displacement and velocity are assumed zero. Forty beam elements of the same length are used for F.E. discretization. To simulate asynchronous support excitation that induces a forced deformation, it is assumed that the seismic acceleration in **Figure 5(c)** is applied on the left support (*i.e.*, support 1) at $t = 0.0$ s and that the earthquake traveling wave propagates longitudinally from the left support to the right ones at constant speed. Assume that the time delay between supports is 0.1 s. For the integration of the finite element equation of motion, the Newmark integration scheme is employed and the time interval is 1/1000 s.

The analytic series solutions for displacement, slope, acceleration, moment and shear force are obtained by eigenfunction expansion method with 10 modes. The numerical results such as displacement, velocity and acceleration at $x = 30$ m are compared with their analytical solutions in **Figure 6** and **Figure 7**. The FE solutions of displacement, slope, bending moment and shear force along the beams at some instants are also compared with their analytic solutions in **Figure 8** and **Figure 9**. They show that the numerical results agree with analytical responses.

5. Conclusion

FE dynamic analysis or time history analysis on the two-span Rayleigh-damped Bernoulli-Euler beams subjected to asynchronous support motions is carried out by using the moving support element. And the corresponding analytical solutions are obtained by using eigenfunction expansion method with 10 modes. The numerical results such as displacement, velocity, acceleration, slope, bending moment and shear force are compared

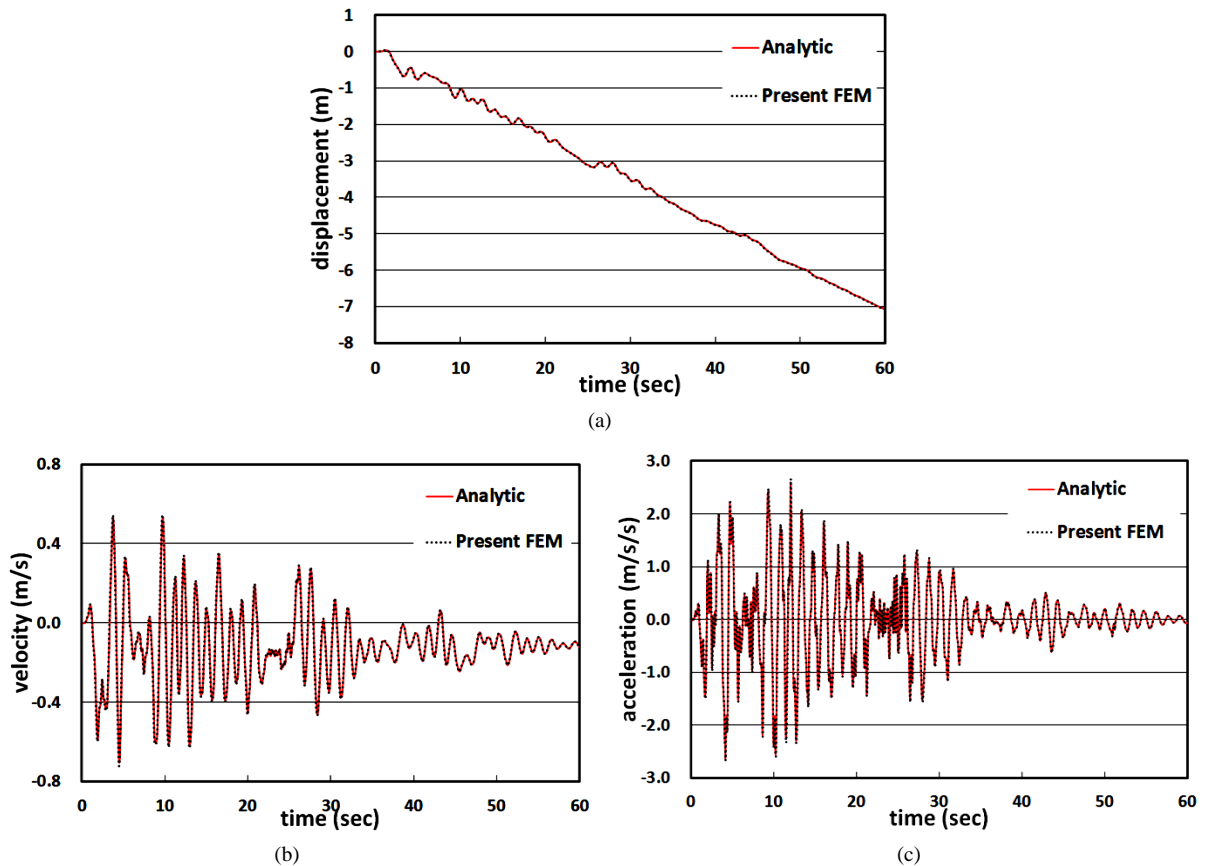


Figure 6. The motions of the beam in **Figure 5(a)**. (a) Displacement at $x = 30$ m; (b) Velocity at $x = 30$ m; (c) Acceleration at $x = 30$ m.

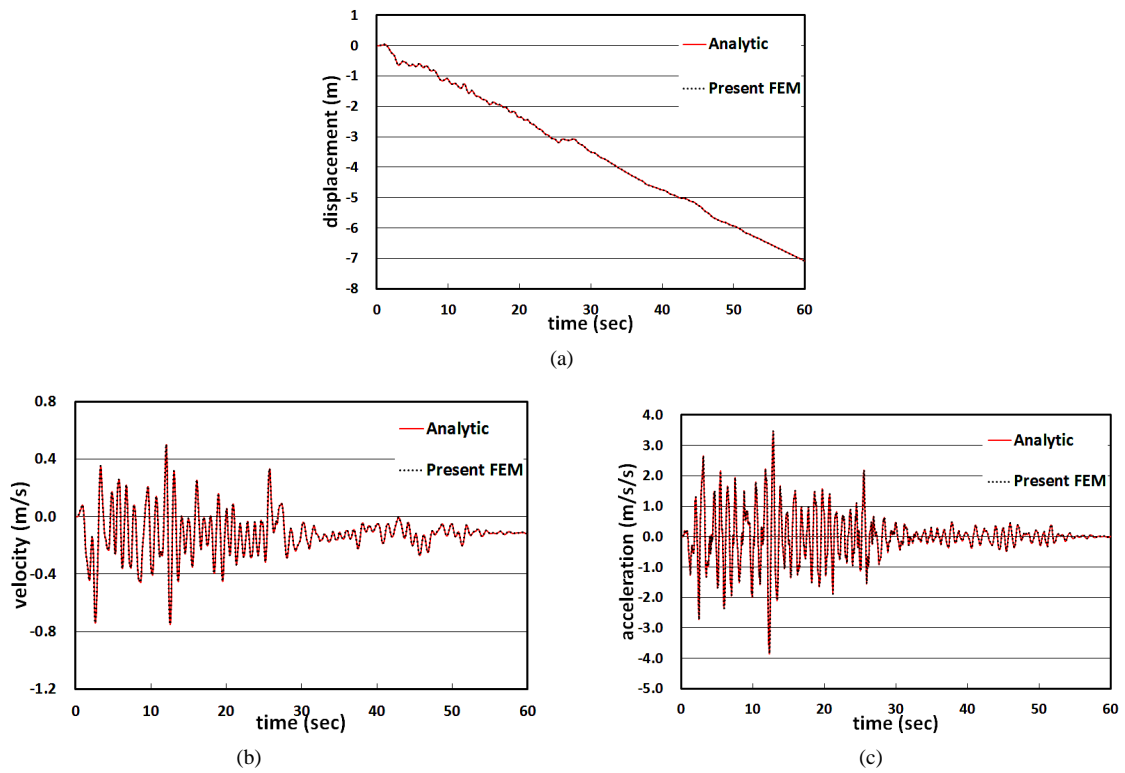


Figure 7. The motions of the beam in **Figure 5(b)**. (a) Displacement at $x = 30$ m; (b) Velocity at $x = 30$ m; (c) Acceleration at $x = 30$ m.

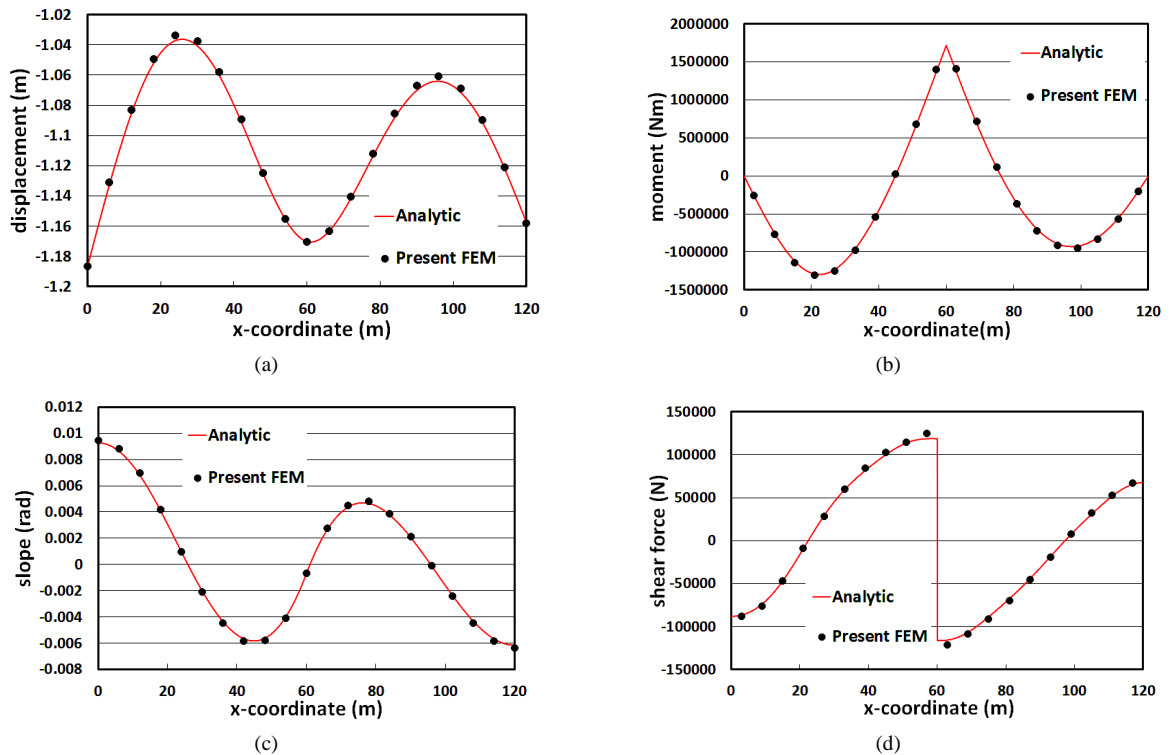


Figure 8. Responses of the beam in **Figure 5(a)** along the beam at $t = 10.18$ sec. (a) Displacement along the beam; (b) Moment along the beam; (c) Slope along the beam; (d) Shear force along the beam.

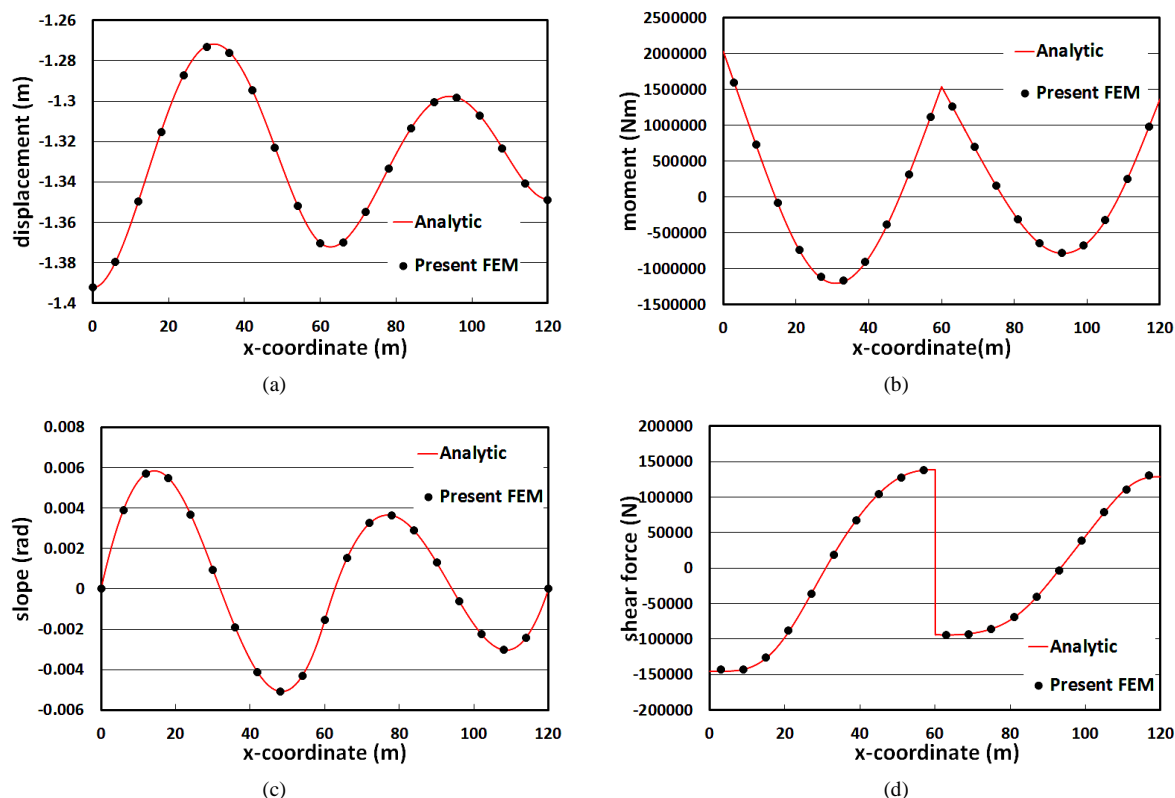


Figure 9. Responses of the beam in Figure 5(b) along the beam at $t = 12.34$ sec. (a) Displacement along the beam; (b) Moment along the beam; (c) Slope along the beam; (d) Shear force along the beam.

with the analytical ones to show that the moving support element describes moving, deforming and vibrating of multi-span beams subjected to support motions accurately. The numerical results agree with analytical solutions well.

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